The motion of billiards in ellipses

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The optical property of ellipses is well known, and also the equivalence: equal angles $\iff \overline{F_1P} + \overline{F_2P} = \text{const.} \iff P \in e$.

There is a generalization:





If any ray is reflected in a conic e then the incoming and the outgoing ray are tangent to the same confocal conic c, called **caustic**.



Charles **Graves** (1812-1899), bishop of Limerick and mathematician:

The locus of point *P* used to pull the string taut around *c* is a confocal ellipse *e*. $D_e := \overline{PQ_1} + \overline{PQ_2} - \overline{Q_1Q_2} = \text{const.}$





Billiards in an ellipse *e* are always tangent to a confocal ellipse or hyperbola.

If one billiard closes after N reflections, then all billiards close, independent of the initial point on c (**Poncelet porism**), and all these closed loops have the same length.



For centuries, billiards attracted the attention of mathematicians, beginning with J.-V. Poncelet and A. Cayley.

S. Tabachnikov: *Geometry and Billiards.* American Mathematical Society, 2005

Recently, Dan Reznik revitalized the interest by computer animations showing the variation of periodic billiards. He identified 40 invariants, e.g., a constant sum of Cosines of interior angles.

Sum of square altitudes to N-periodic tangents is invariant The two types of self-intersected 7-periodics in the Elliptic Billiard





A family of confocal central conics

$$\frac{x^2}{a^2+k} + \frac{y^2}{b^2+k} = 1,$$

 $k \in \mathbb{R} \setminus \{-a^2, -b^2\}$ sends through each point *P* outside the axes one ellipse and one orthogonally intersecting hyperbola.

The parameters (k_e, k_h) define the *elliptic coordinates* of *P* with

 $-a^2 < k_h < -b^2 < k_e$.



We specify the caustic *c* (semiaxes a_c , b_c) as k = 0 and the ellipse *e* with semiaxes a_e , b_e as $k = k_e \implies k_e = a_e^2 - a_c^2 = b_e^2 - b_c^2 > 0$.



From given
$$(k_e, k_h)$$
 follows

$$x^2 = \frac{(a^2 + k_e)(a^2 + k_h)}{d^2},$$

$$y^2 = -\frac{(b^2 + k_e)(b^2 + k_h)}{d^2}.$$
Conversely, $P = (a_e \cos t, b_e \sin t)$
with tangent vector

$$\mathbf{t}_e(t) = (-a_e \sin t, b_e \cos t)$$
and $\mathbf{t}_c(t) = (-a_c \sin t, b_c \cos t)$ yield

$$k_h(t) = -(a_c^2 \sin^2 t + b_c^2 \cos^2 t) = -\|\mathbf{t}_c(t)\|^2 = -\|\mathbf{t}_e(t)\|^2 + k_e.$$

Points on confocal ellipses e and c with the same parameter t have the same coordinate k_h , i.e., they belong to the same confocal hyperbola.



Given $P = (a_e \cos t, b_e \sin t)$, the tangents from P to $c \ (k = 0)$ include the angle $\theta(t)$ where

$$\sin^2 \frac{\theta(t)}{2} = \frac{k_e}{\|\mathbf{t}_e(t)\|^2},$$
$$\tan \frac{\theta(t)}{2} = \pm \frac{\sqrt{k_e}}{\|\mathbf{t}_c(t)\|},$$
$$\cos \theta = 1 - \frac{2k_e}{\|\mathbf{t}_e(t)\|^2} = \frac{k_h(t) + k_e}{k_h(t) - k_e}.$$





The signed distances of the side $t_Q = P_1P_2$ from the center *O* and the pole R_1 w.r.t. *e* have the constant product $-k_e$.

For
$$Q_1 = (a_c \cos t'_1, b_c \sin t'_1),$$

 $\overline{Ot_Q} = \frac{-a_c b_c}{\|\mathbf{t}_c(t'_1)\|}$
 $\overline{Rt_Q} = \overline{R_1 Q_1} = \frac{k_e \|\mathbf{t}_c(t'_1)\|}{a_c b_c}.$





The side P_1P_2 (parameters t_1, t_2) contacts the caustic c iff $\sin^2 \frac{t_1 - t_2}{2} = \frac{k_e}{a_e b_e} \left\| \mathbf{t}_e \left(\frac{t_1 + t_2}{2} \right) \right\|^2$. If t'_1 is the parameter of the tangency point $Q_1 \in c$, then $\tan t'_1 = \frac{b_c a_e}{a_c b_e} \tan \frac{t_1 + t_2}{2}$.

Half angle substitution yields ...



The half-angle substitution $\tau_i := \tan \frac{t_i}{2}$ yields for P_1 and P_2 a symmetric biquadratic equation in projective coordinates on e, namely

 $b_e^2 k_e \tau_1^2 \tau_2^2 - b_c^2 a_e^2 (\tau_1^2 + \tau_2^2) + 2(a_e^2 k_e + a_c^2 b_e^2) \tau_1 \tau_2 + b_e^2 k_e = 0,$

which defines a 2-2-correspondence between consecute points P_1 , P_2 of a billiard. The same holds after N iteration between P_1 and P_{N+1} .

We recall a classical algebraic argument for the Poncelet porism:

A 2-2-correspondence (\neq id) has at most four fixed points. However, four fixed points are already known as contact points between *e* and the common (isotropic) tangents with the caustic *c*. Hence, if one *N*-sided billiard closes, then the correspondence is the identity, and each billiard inscribed in *e* with caustic *c* must close.





S. Tabachnikov: the key result for the integrability of billiards is the **Joachimsthal** integral $J_e := -\langle \mathbf{u}_i, \mathbf{n}_i \rangle$ with \mathbf{u}_i as unit vector of $P_i P_{i+1}$ and $\mathbf{n}_i = (\cos t/a_e, \sin t/b_e)$ as a normal vector of e. This holds in all dimensions $(\mathbf{n}_i = \mathbf{A} \mathbf{p}_i = \mathbf{p}_i^*)$. In the plane

$$J_e = -\langle \mathbf{u}_i, \mathbf{n}_i \rangle = -\cos\left(\frac{\pi}{2} + \frac{\theta_i}{2}\right) \|\mathbf{n}_i\| = \sin\frac{\theta_i}{2} \|\mathbf{n}_i\| = \frac{\sqrt{k_e}}{\|\mathbf{t}_e\|} \frac{\|\mathbf{t}_e\|}{a_e b_e} = \frac{\sqrt{k_e}}{a_e b_e}$$



The extended sides of a billiard intersect at points of confocal ellipses and hyperbolas and form a **Poncelet grid**.

affinely transformed 72-sided periodic billiard with associated Poncelet grid (G. Glaeser, B. Odehnal, H.S.: *The Universe of Conics*, 7KB!)





Theorem:

Given a quadrilateral t_1, \ldots, t_4 of tangents to c from $A_1, B_1 \in c_1$.

Then the range \mathcal{R}_c spanned by c and c_1 contains conics c_2 , c_3 passing through the remaining pairs of opposite vertices (A_2, B_2) and (A_3, B_3) .

(**range** = 'dual pencil')

= summary of results from M. Chasles (1843), W. Böhm (1961), Izmestiev & Tabachnikov (2016), Akopyan & Bobenko (2017).



 $\mathcal{R}_{c} \cap \mathcal{R}_{t} = \{c\} \implies \text{they}$ span a net \mathcal{N} (2-parameter set).

In \mathcal{N} , the line elements of c_1 at A_1 and B_1 span a range which contains the rank-1 conic R.







In \mathcal{N} , the line pencils A_i , B_i and the pencil R (2-fold) span a range which intersects \mathcal{R}_c at c_i . The range contains conics sharing the line elements at A_i and B_i .

The tangents to c_i at A_i and B_i pass through R.

Confocal $c, c_1 \implies$ concyclic quadrilateral.

This holds also when $B_2 \in c$ $(t_1 = t_2)$.













 $D_{e|1} := \overline{Q_1 S_2^{(1)}} + \overline{S_2^{(1)} Q_3} - Q_1 Q_3 = (r_2 + l_2 + w) + (r_3 + l_3 + w) - Q_1 Q_3 = 2D_e + 2w.$



The same invariants show up on the sphere. All circular arcs in black have the same length. The same 'in green'.

In the plane: w =

$$=\frac{2a_eb_e\sqrt{k_e^3}}{a_c^2b_c^2-k_e^2}$$





For each billiard $P_0P_1P_2...$ in the ellipse *e* with caustic *c*, there exists a **conjugate** billiard $P'_0P'_1P'_2...$

The axial scaling with $c \rightarrow e$,

 $lpha\colon Q_i\mapsto P_i',\quad Q_{i-1}'\mapsto P_i$,

transforms tangents P_1P_2 of c to tangents $P'_1P'_2$. This results from the symmetry between t_i and t'_i in the equation

 $b_c a_e \cos t_i \cos t'_i + a_c b_e \sin t_i \sin t'_i = a_c b_c$, which expresses that P_i lies on the tangent to c at Q_i and P'_i on the tangent at Q'_i .





One might say, point P'_i is halfway from P_i to P_{i+1} .

In the sequence of parameters t_1 , t'_1 , t_2 , t'_2 , ... for P_1 , $P'_1(Q_1)$, P_2 , $P'_2(Q_2)$, ..., the transition $P_i \mapsto P'_i$ means a shift $t_i \mapsto t'_i$.

I. Izmestiev, S. Tabachnikov (2017): There exists a **canonical parametrization** of *e* such that the billiard transformation $P_i \rightarrow P_{i+1}$ corresponds to a shift $u_i \rightarrow u_{i+1} = u_i + 2\Delta u$.

Above, an example of canonical parameters: $P_1 \sim 0$, $P_2 \sim \frac{1}{2}$, $P_3 \sim 1$.



3. Periodic N-sided billiards



The **turning number** τ of a periodic billiard counts the loops around the center.

Theorem:

(i) A periodic billiard with even N and odd τ is centrally symmetric.

(ii) For odd N and odd τ , the billiard is centrally symmetric to the conjugate billiard.

(iii) For odd N and even τ the billiard conincides with the conjugate billiard.



3. Periodic N-sided billiards



From Graves' theorem

$$D_e = \overline{Q_{i-1}P_i} + \overline{P_iQ_i} - Q_{i-1}Q_i$$

follows for the perimeter of the *N*-sided billiard

 $L_e = N \cdot D_e + \tau \cdot P_c$

with P_c as perimeter of e.

Ivory's theorem implies for **odd** N = 2n + 1: the length l_i equals symmetric l'_{i+n} of the conjugate billiard and r_{i+n+1} of the original one.

Theorem: $\sum l_i = \sum r_i = L_e/2.$



3. Periodic N-sided billiards



Theorem [Akopyan, Schwartz, Tabachnikov, Bialy] $\sum_{i=1}^{N} \cos \theta_i = N - \frac{\sqrt{k_e}}{a_e b_e} L_e$.

With
$$\sum_{i=1}^{N} \frac{1}{\|\mathbf{t}_e(t_i)\|^2}$$
 also $\sum_{i=1}^{N} \overline{Ot_P}^2$ and $\sum_{i=1}^{N} \kappa_e(t_i)^{2/3}$ are invariant.





With a **billiard motion** we denote a variation of the billiard and the Poncelet grid induced by the variation of a single vertex.

According to Graves' construction, the \mathbf{v}_2 of P_2 can be decomposed as

 $\mathbf{v}_2 = \mathbf{v}_{t_1} + \mathbf{v}_{n_1} = \mathbf{v}_{t_2} + \mathbf{v}_{n_2},$ where due to the constant length $\|\mathbf{v}_{t_2}\| = \|\mathbf{v}_{t_1}\|$ and $\|\mathbf{v}_{n_2}\| = \|\mathbf{v}_{n_1}\|$ $l_2 \omega_2 = r_2 \omega_1,$ hence $\frac{\omega_1}{\omega_2} = \frac{l_2}{r_2}.$

If the billiard is periodic, then the product of all ratios I_i/r_i yields

$$\frac{l_1}{r_1} \cdot \frac{l_2}{r_2} \cdots \frac{l_N}{r_N} = \frac{\omega_N}{\omega_1} \cdot \frac{\omega_1}{\omega_2} \cdots \frac{\omega_{N-1}}{\omega_N} = 1, \text{ hence } l_1 l_2 \dots l_N = r_1 r_2 \dots r_N$$



For periodic billiards, a given velocity vector \mathbf{v}_2 of any vertex defines all velocities.

In terms of the exterior angles $\theta_1, \ldots, \theta_N$ we obtain $\sin \frac{\theta_2}{2} = \frac{l_2 \omega_2}{v_2} = \frac{r_2 \omega_1}{v_2}$ and $\cos \frac{\theta_2}{2} = \frac{v_{t|2}}{v_2}$, where $v_2 := ||\mathbf{v}_2||, v_{t|2} := ||\mathbf{v}_{t|2}||.$ From $P_1 P_2 \perp Q_1 R_1$ follows $\overline{R_1 Q_1} = l_1 \tan \frac{\theta_1}{2} = r_2 \tan \frac{\theta_2}{2}.$





After some manipulations follows $v_{t|1} \tan^2 \frac{\theta_1}{2} = v_{t|2} \tan^2 \frac{\theta_2}{2}$ $\dots = v_{t|i} \frac{k_e}{\|\mathbf{t}_c(t_i)\|^2} =: C.$ Instead of a free choice of v_2 , we set $C = k_e$.

 $v_{t|i} = \|\mathbf{t}_{c}\|^{2}, \quad v_{i} = \|\mathbf{t}_{c}\| \|\mathbf{t}_{e}\|,$ $v_{n|i} = v_{i} \sin \frac{\theta}{2} = \|\mathbf{t}_{c}\| \sqrt{k_{e}}$ for all $t = t_{i}$ and for all confocal ellipses e.





To each point $\mathbf{p} = (x, y) = (a_e \cos t, b_e \sin t)$ we assign a velocity vector $\mathbf{v} = \|\mathbf{t}_c\| \mathbf{t}_e = \sqrt{a_c^2 \sin^2 t + b_c^2 \cos^2 t} \left(-\frac{a_e y}{b_e}, \frac{b_e x}{a_e}\right).$

Theorem:

This vector field defines an infinitesimal motion which preserves confocal ellipses and permutes the confocal hyperbolas and the tangents of the caustic *c*.

The infinitesimal motion generates a one-parameter Liegroup Γ which carries out the billiard transformation along e and simultaneously that of the associated Poncelet grid.





We set $\mathbf{v} = \frac{\mathrm{d} \mathbf{p}}{\mathrm{d} u} = \mathbf{p}$. Then, u is a canonical parameter of e and of Γ , i.e., $\gamma(u_2) \circ \gamma(u_1) = \gamma(u_1 + u_2)$. $\mathbf{v}(t) = \|\mathbf{t}_c(t)\| \mathbf{t}_e(t) = t \mathbf{t}_e(t)$ $\dot{t} = \frac{\mathrm{d} t}{\mathrm{d} u} = \sqrt{a_c^2 \sin^2 t + b_c^2 \cos^2 t}$.

Proof: Γ permutes hyperbolas since \dot{t} (and \dot{k}_h) is independent of e (and k_e).

The condition $b_c a_e \cos t \cos t' + a_c b_e \sin t \sin t' = a_c b_c$ is equivalent to the fact that $P = (a_e \cos t, b_e \sin t)$ lies on the tangent t_Q of $Q = (a_c \cos t' + b_c \sin t')$. Differentiation by u yields an identity. Hence, Γ permutes the tangents of c.

In order to express the action of $\gamma(u) \in \Gamma$ on $(a_e \cos t, b_e \sin t)$, we integrate

$$\frac{\mathrm{d}t}{\mathrm{d}u} = \sqrt{a_c^2 \sin^2 t + b_c^2 \cos^2 t} = \sqrt{a_c^2 \sin^2 t + (a_c^2 - d^2) \cos^2 t} = a_c \sqrt{1 - m^2 \cos^2 t}$$

with $m := d/a_c < 1$ as numeric eccentricity of the caustic *c*. We substitute $\varphi := t - \frac{\pi}{2}$ and get under the initial condition $\varphi = 0$ for u = 0

$$\frac{\mathrm{d}\varphi}{\sqrt{1-m^2\sin^2\varphi}} = a_c\,\mathrm{d}u\,, \text{ hence } a_c\,u(\varphi) = F(\varphi,\,m) = \int_0^\varphi \frac{\mathrm{d}\varphi}{\sqrt{1-m^2\sin^2\varphi}}$$

with $F(\varphi, m)$ as the elliptic integral of the first kind with the modulus m. This function shows the canonical coordinate u in terms of φ with the quarter period

$$K := a_c u\left(\frac{\pi}{2}\right) = \int_0^{\pi/2} \frac{\mathrm{d}\varphi}{\sqrt{1 - m^2 \sin^2 \varphi}}$$



For the sake of simplicity, we define $\tilde{u}(\varphi) := a_c u(\varphi)$ as a new canonical coordinate.

The inverse function of $\tilde{u} = F(\varphi, m)$, the Jacobian amplitude $\varphi = \operatorname{am}(\tilde{u})$ leads to the Jacobian elliptic functions,

 $\operatorname{sn} \widetilde{u} = \operatorname{sin}(\operatorname{am}(\widetilde{u}))$ and $\operatorname{cn} \widetilde{u} = \operatorname{cos}(\operatorname{am}(\widetilde{u}))$,

which can be extended in \mathbb{R} to periodic functions with period 4K.



Theorem:

For the ellipse c with semiaxes (a_c, b_c) and eccentricity $d = \sqrt{a_c^2 - b_c^2}$ and all confocal ellipses e with semiaxes (a_e, b_e) , the inscribed billiards with caustic c can be canonically parametrized as $(-a_e \operatorname{sn} \widetilde{u}, b_e \operatorname{cn} \widetilde{u})$, using the Jacobian elliptic functions to the modulus $m = d/a_c$.

If $b_c = b_e \operatorname{cn}(\Delta \widetilde{u})$, then the vertices of the billiard in *e* have the canonical parameters $\widetilde{u} = (\widetilde{u}_0 + 2k\Delta \widetilde{u})$ for $k \in \mathbb{Z}$ and any given initial \widetilde{u}_0 .

Conversely, we obtain an *N*-sided billiard with turning number τ , where $gcd(N, \tau) = 1$, by the choice $\Delta \widetilde{u} = \frac{2\tau K}{N}$ with *K* as the complete elliptic integral of the first kind to the modulus *m*, provided that $b_e = b_c/cn(\Delta \widetilde{u})$.







Schönbrunn Castle, Vienna

Thank you for your attention!



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